Lippman-Schwinger-Lanczos algorithm for inverse scattering problems.

¹WPI

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Background

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- Forward PDE problems: Given the PDE, including its coefficients, find its solution everywhere.
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- Reduced Order Models (ROMs) for forward problems: If e.g. PDE is linear, find a low dimensional matrix that acts like the differential operator.
- ROMs for inverse problems: Given data, find a ROM which matches the data, use this ROM to extract the unknown coefficient.

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- (Druskin, Zaslavsky, M 2021) Use data generated internal solution in a Lippmann-Schwinger formulation.
- Time domain reconstruction of wave speed (Borcea, Garnier, Mamonov, Zimmerling 2022)

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- Druskin, V., Mamonov, A. and Zaslavsky, M., A nonlinear method for imaging with acoustic waves via reduced order model backprojection, SIAM Journal on Imaging Sciences, (2018).
- Borcea, L., Druskin, V., and Mamonov, A., Zaslavsky, M. and Zimmerling, J., Reduced Order Model Approach to Inverse Scattering, SIAM Journal on Imaging Sciences, (2020).

$$u_{tt} + Au = 0 \text{ in } \Omega \times [0, \infty) \tag{1}$$

$$u(t=0) = g \text{ in } \Omega \tag{2}$$

$$u_t(t=0) = 0 \text{ in } \Omega \tag{3}$$

where

$$A = A_0 + q \tag{4}$$

- $A_0 \ge 0$ is known background, (for example $A_0 = -\Delta$),
- $q(x) \ge 0$ is our unknown potential
- initial data g is localized (approximate delta) source
- assume homogeneous Neumann boundary conditions on the spatial boundary ∂Ω.

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$$u(x,t) = \cos\left(\sqrt{A}t\right)g(x). \tag{5}$$

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• The inverse problem is as follows: Given

$$\{F(k\tau)\}\$$
 for $k = 0, \dots, 2n-2,$

reconstruct q.

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- then the $n \times n$ mass matrix $k, l = 0, \ldots, n-1$

$$M_{kl} = \int_{\Omega} u_k u_l dx \tag{7}$$

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• from the cosine angle sum formula

$$M_{kl} = \frac{1}{2} \left(F((k-l)\tau) + F((k+l)\tau) \right),$$
(9)

M can be obtained directly from the data.

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$$M = U^{\top}U$$

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 The functions {v_k} will be orthonormal in the L² norm (Gram-Schmidt). • We do not know the snapshots, but from the data we know the transformation that orthogonalizes them sequentially.

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- So do all of the above for the known background problem

Background exact solution

$$u^{0}(x,t) = \cos(\sqrt{A_{0}}t)g(x).$$
 (11)

and snapshots $\{u_i^0\}$

mass matrix

$$M_{kl}^0 = \int_{\Omega} u_k^0 u_l^0 dx, \qquad (12)$$

Cholesky decomposition

$$M^0 = (U^0)^\top U^0,$$

orthogonalized background snapshots

$$\vec{v}^0 = \vec{u}^0 (U^0)^{-1}.$$
 (13)

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• Definition of our data generated snapshots

$$\vec{\mathbf{u}} := \vec{v}^0 U = \vec{u}^0 (U^0)^{-1} U.$$
(15)



Figure: Data generated internal snapshots

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• Time domain Lippmann-Schwinger

$$F_0(k\tau) - F(k\tau) = \int_0^{k\tau} \int_{\Omega} u_0(x, k\tau - t) u(x, t) q(x) dx dt.$$
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• Use data generated internal solution (interpolated in time)

$$F_0(k\tau) - F(k\tau) = \int_0^{k\tau} \int_{\Omega} u_0(x, k\tau - t) \mathbf{u}(x, t) q(x) dx dt \qquad (17)$$

• Given u such that

$$-u'' + q(x)u + \lambda u = 0$$
 for x on (0,1)
 $-u'(0) = 1$
 $u(1) = 0$

- Define the transfer function $F(\lambda) := u(0; \lambda)$.
- Consider the inverse problem: Given $\{F(\lambda), F'(\lambda) : \lambda = b_1, \dots b_m\}$, find q(x)

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- Given 2m spectral data values to reconstruct q(x)
- Can do a modified version of what follows for other forms of spectral data
- We will construct a ROM that matches this data exactly

Spectral domain SISO

 Consider exact solutions to above u₁,..., u_m corresponding to spectral points λ = b₁,... b_m. and the subspace

$$G = \operatorname{span}\{u_1, \ldots, u_m\}$$

- Although we do not know these solutions, we can obtain the Galerkin system (ROM) from the data
- Given by the mass and stiffness matrices

$$M_{ij}=\int_0^1 u_i u_j$$

and

$$S_{ij} = \int_0^1 u'_i u'_j + \int_0^1 q u_i u_j.$$

They are given by the formulas

$$M_{ij} = \frac{F(\lambda_i) - F(\lambda_j)}{\lambda_j - \lambda_i}, \quad M_{ii} = -\frac{dF}{d\lambda}(\lambda_i).$$
(18)

and

$$S_{ij} = \frac{F(\lambda_j)\lambda_j - F(\lambda_i)\lambda_i}{\lambda_j - \lambda_i}, \quad S_{ii} = \frac{d(\lambda F)}{d\lambda}(\lambda_i).$$
(19)

• Searching for the unknown coefficients $\{c_i\}$ for the solution

$$u_G = \sum_{i=1}^m c_i u_i$$
$$M_{ij} = \int_0^1 u_i u_j$$

and

$$S_{ij} = \int_0^1 u'_i u'_j + \int_0^1 q u_i u_j.$$

• For forward solution would solve $(S + \lambda M)\vec{c} = \vec{F}$ where $F_i = F(b_i) = u_i(0)$.

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- Why does this work in the spectral domain?
- because the Krylov subspaces *are the same as* those generated by ROM- projected sequential time snapshots.

Spectral domain SISO

• That is, if $d \in \mathbb{R}^m$ satisfies the Galerkin problem

$$Sd(t) + Md(t)_{tt} = 0, \qquad d(0) = b, \qquad d_{tt=0} = 0,$$

which is a time-domain (the wave) variant of the ROM.

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• Then $d(\tau i)$ satisfy the second order finite-difference scheme

$$d[\tau(i+1)] = (2I - \tau A)d[\tau i] - d[\tau(i-1)], i = i, \dots, m-1,$$

$$d(0) = M^{-1}b, \quad d(\tau) = d(-\tau)$$

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 span{d(\(\tau\)i)} are the same as the above Krylov subspaces w/ powers of A. • So the entries of this orthogonalized reduced order model (which can be obtained from the data) are the entries of the stiffness matrix

$$\hat{S}_{ij}=\int \hat{u}_i'\hat{u}_j'+\int_0^1 q\hat{u}_i\hat{u}_j$$

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• correspond to sequentially orthogonalized projected time snapshots, which *depend only very weakly on the coefficient*.

Weak dependence of orthogonalized spectral basis on q (positive $\lambda = b_i$)



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reference medium

Spectral domain Lippmann-Schwinger Lanczos approach

• Since the basis depends weakly on the unknown coefficient, we can get a data generated internal solution

$$\mathbf{u} = \sqrt{b^* M^{-1} b} V_0 Q_0 (A + \lambda I)^{-1} e_1$$

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$$F_0 - F = \int_{\Omega} u u_0 (q - q_0) \tag{20}$$

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- For inverse Born one would replace u by u_0
- With data generated ROM we replace by **u** , the data generated internal solution.

Spectral domain Lippman-Schwinger Lanczos approach







Spectral domain Lippman-Schwinger Lanczos approach



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L-S-L for inverse scattering

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Spectral domain Lippman-Schwinger Lanczos approach



Figure: Experiment 3: True medium (top left) and its reconstructions using 'Cheated IE' (top right), Born linearization (bottom left) and our approach (bottom right)

symmetric data: Lippman-Schwinger Lanczos approach



Figure: Experiment 1: True medium (top left) and its reconstructions using 'Cheated IE' (top right), Born linearization (bottom left) and our approach (bottom right)

Non-symmetric data: Lippman-Schwinger Lanczos approach



Figure: Experiment 1: True medium (top left) and its reconstructions using 'Cheated IE' (top right), Born linearization (bottom left) and our approach (bottom right)

Helmholtz 2d (with E. Cherkaev and J. Baker)



Figure: 2-D Helmholtz (positivé 3)

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L-S-L for inverse scattering

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- In both cases the new basis is close to that from reference medium.
- Can use the reference medium basis to obtain approximations of internal solutions from data only
- Lippmann-Schwinger-Lanczos: use these internal solutions in Lippmann Schwinger, extendable to more general data sets.